

# Modifying gravity through instabilities

Fethi M Ramazanoğlu  
Koç University

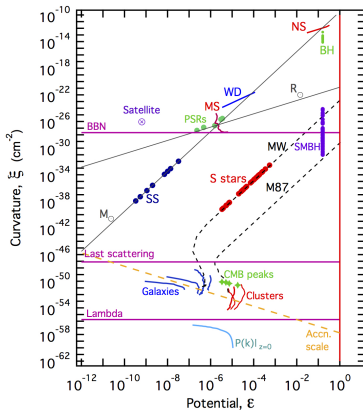
“GWverse: black holes, gravitational waves and fundamental physics”  
Supported under CA 16104

University of Malta  
January 24, 2018

# Takeaway Message

- Spontaneous scalarization is a member of a family of theories common in their underlying mechanism and observable signatures: **spontaneous tensorization**.
- They all arise from a spontaneously growing instabilities which are eventually regularized: **regularized instabilities**.
- They all (likely) have order-of-unity deviations from GR in strong field: **relevant for near-future observations**.

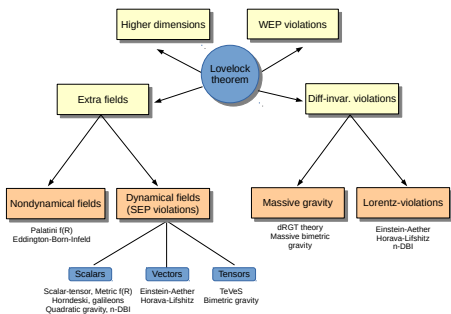
# Testing general relativity



Baker et al. 2014

- Many since Eddington's observation of light deflection in 1919
- GR is victorious, including the “bizarre” predictions
- So far mainly for **weak fields**.
- Field theories most likely break at strong-fields.
- Now we have the means to probe **dynamical strong fields**, even dynamical ones: **gravitational waves**.

# Modifying general relativity



Berti et al. 2015

- Many motivations from unification, high energy ...
- Weak field well tested: cannot modify much (killed many early ideas)
- Strong field experimental accuracy low: large modifications preferred
- Make sure there is no fundamental physical or mathematical problem
- Still left with many options

# Scalar-Tensor theories & Spontaneous Scalarization

$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} R \\ + S_M(g_{ab}, \psi)$$

# Scalar-Tensor theories & Spontaneous Scalarization

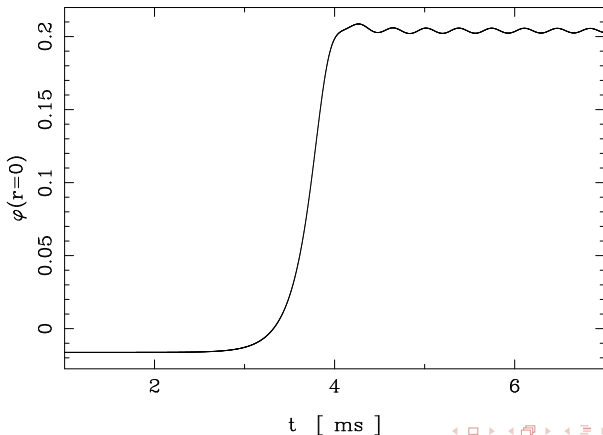
$$S = \frac{1}{16\pi G} \int d^4x \sqrt{-g} R - \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left( 2(\partial\phi)^2 + 2m_\phi\phi^2 \right) + S_M(A^2(\phi)g_{ab}, \psi)$$

$$\tilde{g}_{ab} \equiv A^2(\phi)g_{ab}$$

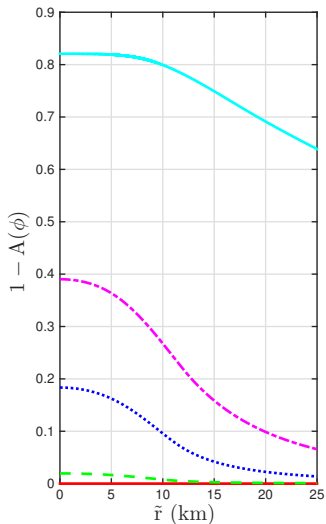
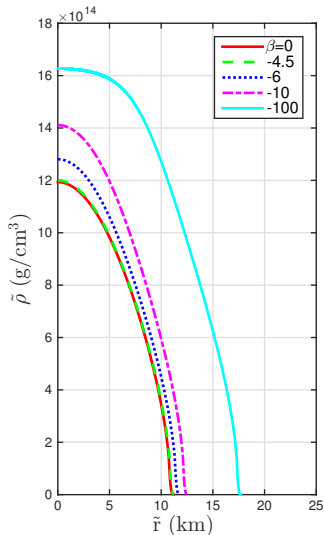
- Possibly most famous alternative to GR.
- Less popular after solar system observations
- Many  $f(R)$  theories can be recast as scalar-tensor theories.
- Spontaneous scalarization when  $A(\phi) = e^{\beta/2} \phi^2$  (Damour, Esposito-Farese, 1993)

# Why Spontaneous?

Novak, 1998



# How does this modification effect neutron stars?





# Why Spontaneous?: Tachyon Instability

$$\begin{aligned}\square_g \phi &= \left( m_\phi^2 - 8\pi A^4 \frac{d \ln A}{d(\phi^2)} \right) \phi \\ &\approx \left( m_\phi^2 - 4\pi \beta \tilde{T} \right) \phi \\ &\approx m_{\text{eff}}^2 \phi\end{aligned}$$

- Non-relativistic Matter:  $T = -\rho + 3P \approx -\rho < 0$ .
- $m_{\text{eff}}^2 < 0$ , tachyon instability!
- Alternatively,  $V_{\text{eff}}$  is not bounded from below.
- Not tachyonic for larger  $\phi$ : regularization.

# Generalizing to other fields

$$S = \underbrace{\frac{1}{16\pi G} \int d^4x \sqrt{-g} R}_{S_{GR}} - \underbrace{\frac{1}{16\pi G} \int d^4x \sqrt{-g} \left( 2(\partial\phi)^2 + 2m_\phi \phi^2 \right)}_{S_\phi} + \underbrace{S_M(A^2(\phi)g_{ab}, \psi)}_{S_{matter}}$$

Recipe: Replace  $\phi$  with any other field and choose  $A$  as an “inverse parabola”. The new field will have a tachyonic EOM.

# Example: Spontaneous vectorization

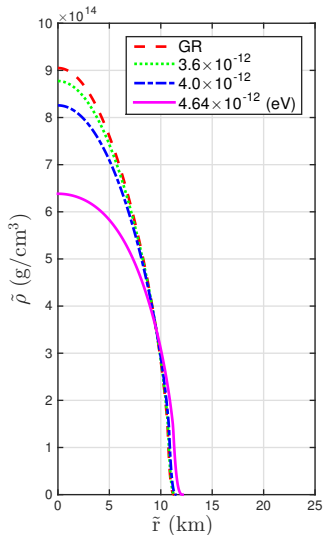
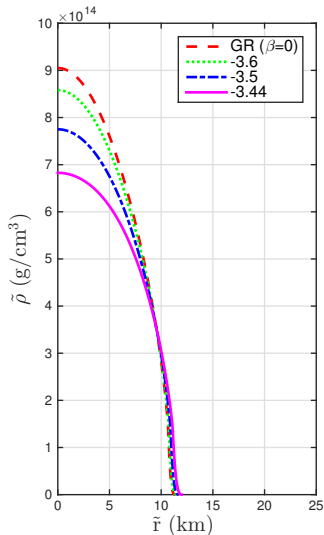
$$S = \underbrace{\frac{1}{16\pi G} \int d^4x \sqrt{-g} R}_{S_{GR}} - \underbrace{\frac{1}{16\pi G} \int d^4x \sqrt{-g} \left( F^{\mu\nu} F_{\mu\nu} + 2m_X^2 X_\mu X^\mu \right)}_{S_{X^\mu}}$$

$$+ \underbrace{S_M(A^2(X_\mu X^\mu) g_{ab}, \psi)}_{S_{matter}} \quad F_{\mu\nu} \equiv \partial_\mu X_\nu - \partial_\nu X_\mu$$

$$\Rightarrow \nabla_\rho F^{\rho\mu} = \left( -8\pi A_X^4 \frac{d \ln A_X}{d(X_\mu X^\mu)} \tilde{T} + m_X^2 \right) X^\mu$$

$A = e^{\beta X_\mu X^\mu}$  again gives negative mass square!

# How does this modification effect neutron stars?



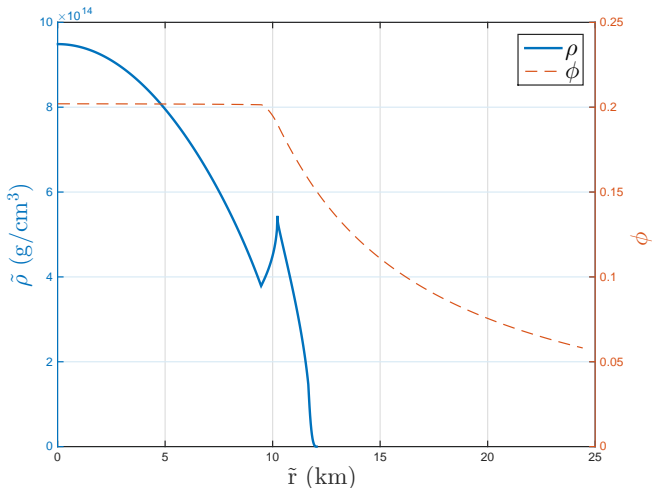
# Generalizing to other instabilities: Ghost-based scalarization

$$S = \underbrace{\frac{1}{16\pi G} \int d^4x \sqrt{-g} R}_{S_{GR}} - \underbrace{\frac{1}{16\pi G} \int d^4x \sqrt{-g} \left( 2(\partial\phi)^2 + 2m_\phi\phi^2 \right)}_{S_\phi} + \underbrace{S_M(A^2(K)g_{ab}, \psi)}_{S_{matter}}, \quad K = (\partial_\mu\phi\partial^\mu\phi)$$

$$\Rightarrow \left( -8\pi A^4 \frac{dA}{dK} \tilde{T} + 1 \right) \square\phi = m_\phi^2\phi$$

Recipe: Replace the potential-like conformal scaling with kinetic-like conformal scaling. The scalar will be a ghost!

# Generalizing to other instabilities: Ghost-based scalarization



# Regularization of instabilities in Gravity

- Can combine the two ideas: **ghost-based spontaneous vectorization**
- Beyond vectors, spin-2 fields? Major obstacles due to “other” ghosts.
- The essence is not scalars vs vectors, or tachyon vs ghost: **instability regularization in GR**
  - Take a field and incite an instability
  - leave it to nonlinearity to regularize it
  - alternative theory with large deviations in strong field
  - ???
  - profit

# Relevance to CO mergers

- Due to large deviations, any mergers with NSs are significant
- Events with mass change are even better
  - NSNS mergers with long-lived HMNS
  - Eccentric mergers with partial disruptions
  - ...



# Summary – Future

- Still quite a bit to do with (massive) scalarization.
- Well-posedness of the family.
- Perturbative signatures: monopole, dipole emissions, ...
- Numerical relativity
- Spin-2?
- Regularized instabilities to any other modifications to GR?